

The Joys of MOIs

University of Bonn
Oberseminar Global Analysis and Operator Algebras

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Summary of this talk

- 1 Motivation for studying MOIs
- 2 Pseudodifferential calculus
- 3 MOIs of pseudodifferential operators

This talk is based on joint work with Edward McDonald and Teun van Nuland.

Part 1: Motivation

Exhibit A

We use the Chern character of $(\mathcal{A}, \mathcal{H}, D)$ in entire cyclic cohomology (cf. [2]) given in the most efficient manner by the JLO formula, which defines the components of an entire cocycle in the (b, B) bicomplex:

$$(90) \quad \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_0^n v_i = 1, v_i \geq 0}$$

$$\text{Trace} \left(a^0 e^{-v_0 D^2} [D, a^1] e^{-v_1 D^2} \dots e^{-v_{n-1} D^2} [D, a^n] e^{-v_n D^2} \right), \quad \forall a^j \in \mathcal{A}$$

where n is odd.

We introduce a parameter ϵ by replacing D^2 by ϵD^2 , which yields a cocycle ψ_n^ϵ which is cohomologous to ψ_n . One has moreover

$$(91) \quad \psi_n^\epsilon(a^0, \dots, a^n) = \sqrt{2i} \left(\int_{\sum_0^n v_i = 1} \theta(\epsilon v_0, \dots, \epsilon v_n) \pi dv_i \right) \epsilon^{n/2},$$

From [ConnesMoscovici1995]

Exhibit B

Let us now show that if $b \in \cap \text{Dom } L^k R^q$ then $b \in \text{Dom } \delta$. The proof is more subtle than one would expect, because the obvious argument, using

$$|D| = \pi^{-1} \int_0^\infty \frac{D^2}{D^2 + \mu} \mu^{-1/2} d\mu ,$$

requires some care. Indeed, one gets from the above

$$[|D|, b] = \pi^{-1} \int_0^\infty (D^2 + \mu)^{-1} [D^2, b] (D^2 + \mu)^{-1} \mu^{1/2} d\mu .$$

We can replace $[D^2, b]$ by $|D|$, which has the same size, and get

$$\int_0^\infty (D^2 + \mu)^{-2} |D| \mu^{1/2} d\mu = \int_0^\infty (1+t)^{-2} t^{1/2} dt .$$

For this to work, we need to move $[D^2, b]$ in front of the above integral, i.e. use the finiteness of the norm of

$$\int_0^\infty \underbrace{[(D^2 + \mu)^{-1}, [D^2, b]]}_{-(D^2 + \mu)^{-1} [D^2, [D^2, b]] (D^2 + \mu)^{-1}} (D^2 + \mu)^{-1} \mu^{1/2} d\mu .$$

This finiteness follows from:

- 1) $(D^2 + \mu)^{-1} [D^2, [D^2, b]]$ bounded since $b \in \text{Dom } L^2$
- 2) $\int_0^\infty \|(D^2 + \mu)^{-2}\| \mu^{1/2} d\mu \leq C \int_0^1 \mu^{1/2} d\mu + \int_1^\infty \mu^{-3/2} d\mu < \infty$.

Once $[D^2, b]$ is moved in front the above calculation applies.

From [ConnesMoscovici1995]

Exhibit C

Now, onwards with the computation, the first part of which is straightforward:

$$\begin{aligned}
 [\Delta^{-z}, A]B &= \frac{1}{2\pi i} \int \lambda^{-z} [(\lambda - \Delta)^{-1}, A]B \, d\lambda \\
 &= \frac{1}{2\pi i} \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] (\lambda - \Delta)^{-1} B \, d\lambda \\
 &= \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] B (\lambda - \Delta)^{-1} \, d\lambda \\
 &\quad + \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] (\lambda - \Delta)^{-1} [\Delta, B] (\lambda - \Delta)^{-1} \, d\lambda.
 \end{aligned}$$

(In the last step we did two things at once: we commuted B past $(\lambda - \Delta)^{-1}$ and we then used the formula $[S^{-1}, T] = S^{-1}[T, S]S^{-1}$.) The operators $[\Delta, A]$ and $[\Delta, B]$ have orders 1 and 2, respectively.

Before going on, we shall introduce some better notation for our contour integrals.

2.5 Definition. If $D_0, \dots, D_p$ are differential operators on the closed manifold M , then denote by $I_z(D_0, \dots, D_p)$ the integral

$$\frac{1}{2\pi i} \int \lambda^{-z} D_0 (\lambda - \Delta)^{-1} \dots D_p (\lambda - \Delta)^{-1} \, d\lambda$$

(in the integral, copies of $(\lambda - \Delta)^{-1}$ alternate with the operators D_j). The integral converges if $\operatorname{Re}(z) < n$, in the sense we discussed above, and defines an operator on $C^\infty(M)$.

From [Higson2003]

Exhibit D

Theorem 4.2 (Semifinite Odd Local Index Theorem). *Let $(\mathcal{A}, \mathcal{H}, \mathcal{D})$ be an odd finitely summable QC^∞ spectral triple with spectral dimension $p \geq 1$. Let $N = [p/2] + 1$ where $[\cdot]$ denotes the integer part, and let $u \in \mathcal{A}$ be unitary. Then*

$$1) \quad sf(\mathcal{D}, u^* \mathcal{D} u) = \frac{1}{\sqrt{2\pi i}} res_{r=(1-p)/2} \left(\sum_{m=1, \text{odd}}^{2N-1} \phi_m^r(Ch_m(u)) \right)$$

where for $a_0, \dots, a_m \in \mathcal{A}$, $l = \{a + iv : v \in \mathbf{R}\}$, $0 < a < 1/2$, $R_s(\lambda) = (\lambda - (1 + s^2 + \mathcal{D}^2))^{-1}$ and $r > 0$ we define $\phi_m^r(a_0, a_1, \dots, a_m)$ to be

$$\frac{-2\sqrt{2\pi i}}{\Gamma((m+1)/2)} \int_0^\infty s^m \tau \left(\frac{1}{2\pi i} \int_l \lambda^{-p/2-r} a_0 R_s(\lambda) [\mathcal{D}, a_1] R_s(\lambda) \cdots [\mathcal{D}, a_m] R_s(\lambda) d\lambda \right) ds.$$

In particular the sum on the right hand side of 1) analytically continues to a deleted neighbourhood of $r = (1-p)/2$ with at worst a simple pole at $r = (1-p)/2$. Moreover, the complex function-valued cochain $(\phi_m^r)_{m=1, \text{odd}}^{2N-1}$ is a (b, B) cocycle for $\mathcal{A}$ modulo functions holomorphic in a half-plane containing $r = (1-p)/2$.

From [CareyPhillipsRennieSukochev2006]

Exhibit E

Let us introduce the following convenient notation (cf. [10]). If $A_0, \dots, A_n$ are operators, we define a t -dependent quantity by

$$\langle A_0, \dots, A_n \rangle_n := t^n \text{Tr} \int_{\Delta_n} A_0 e^{-s_0 t D^2} A_1 e^{-s_1 t D^2} \dots A_n e^{-s_n t D^2} d^n s. \quad (3)$$

Note the difference in notation with [10], for which the same symbol is used for the supertrace of the same expression, rather than the trace. Also, we are integrating over the ‘inflated’ n -simplex $t\Delta^n$, yielding the factor t^n . The forms $\langle A_0, \dots, A_n \rangle$ satisfy, *mutatis mutandis*, the following properties.

Lemma 7. (See [10].) *In each of the following cases, we assume that the operators A_i are such that each term is well defined:*

1. $\langle A_0, \dots, A_n \rangle_n = \langle A_i, \dots, A_n, \dots, A_{i-1} \rangle_n$;
2. $\langle A_0, \dots, A_n \rangle_n = \sum_{i=0}^n \langle 1, \dots, A_i, \dots, A_n, A_0, \dots, A_{i-1} \rangle_n$;
3. $\sum_{i=0}^n \langle A_0, \dots, [D, A_i], \dots, A_n \rangle_n = 0$;
4. $\langle A_0, \dots, [D^2, A_i], \dots, A_n \rangle_n = \langle A_0, \dots, A_{i-1} A_i, \dots, A_n \rangle_{n-1} - \langle A_0, \dots, A_i A_{i+1}, \dots, A_n \rangle_{n-1}$.

From [vanSuijlekom2011]

Multiple operator integrals

Let $\phi : \mathbb{R}^{n+1} \rightarrow \mathbb{C}$ be such that

$$\phi(\lambda_0, \dots, \lambda_n) = \int_{\Omega} a_0(\lambda_0, \omega) \cdots a_n(\lambda_n, \omega) d\nu(\omega),$$

with finite measure space (Ω, ν) and measurable and bounded $a_j : \mathbb{R} \times \Omega \rightarrow \mathbb{C}$.

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with finite measure space (Ω, ν) and measurable and bounded $a_j : \mathbb{R} \times \Omega \rightarrow \mathbb{C}$.

Let $H_0, \dots, H_n$ be self-adjoint, for $V_1, \dots, V_n \in B(\mathcal{H})$ define the MOI

$$\begin{aligned} T_{\phi}^{H_0, \dots, H_n}(V_1, \dots, V_n)\psi \\ := \int_{\Omega} a_0(H_0, \omega) V_1 a_1(H_1, \omega) \cdots V_n a_n(H_n, \omega) \psi d\nu(\omega), \quad \psi \in \mathcal{H}. \end{aligned}$$

Then,

$$T_{\phi}^{H_0, \dots, H_n} : B(\mathcal{H}) \times \cdots \times B(\mathcal{H}) \rightarrow B(\mathcal{H})$$

and this does not depend on how we represent ϕ (its *symbol*).

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Consider for example $f : \mathbb{C} \rightarrow \mathbb{C}$ holomorphic, and $A, B \in B(\mathcal{H})$ self-adjoint. Suppose we want to study $f(A) - f(B)$.

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For a large enough contour γ ,

$$\begin{aligned} f(A) - f(B) &= \frac{1}{2\pi i} \int_{\gamma} f(z)(z - A)^{-1} dz - \frac{1}{2\pi i} \int_{\gamma} f(z)(z - B)^{-1} dz \\ &= \frac{1}{2\pi i} \int_{\gamma} f(z)(z - A)^{-1}(A - B)(z - B)^{-1} dz \\ &= T_{f^{[1]}}^{A,B}(A - B). \end{aligned}$$

The symbol $f^{[1]}$ here is very typical.

Divided differences

Symbols of MOIs encountered in the wild are almost always divided differences, which are defined recursively for $f \in C^n(\mathbb{R})$ as

$$f^{[0]}(\lambda) := f(\lambda);$$

$$f^{[n]}(\lambda_0, \dots, \lambda_n) := \frac{f^{[n-1]}(\lambda_0, \dots, \lambda_{n-1}) - f^{[n-1]}(\lambda_1, \dots, \lambda_n)}{\lambda_0 - \lambda_n},$$

with an appropriate limit if $\lambda_0 = \lambda_n$. In particular,

$$\frac{1}{n!} f^{(n)}(\lambda) = f^{[n]}(\lambda, \dots, \lambda).$$

Example MOIs 2

For example, the JLO cocycle is

$$\begin{aligned} \int_{\Delta_n} \mathrm{Tr}(\eta a_0 e^{-t_0 D^2} [D, a_1] e^{-t_1 D^2} \cdots [D, a_n] e^{-t_n D^2}) dt \\ = \mathrm{Tr}(\eta a_0 T_{f[n]}^{D^2}([D, a_1], \dots, [D, a_n])), \end{aligned}$$

with $f(x) = \exp(-x)$.

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- $\frac{d^n}{dt^n} f(H + tV)|_{t=0} = T_{f^{[n]}}^{H, \dots, H}(V, \dots, V),$

each of which has been used to obtain sharp estimates. (Potapov, Sukochev, Skripka, Caspers, Montgomery-Smith, McDonald, Peller, ...)

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Furthermore, MOIs can systematise operator integral techniques in NCG.

A problem

If you write, like in *The Local Index Formula in Noncommutative Geometry* by Nigel Higson, for a spectral triple $(\mathcal{A}, \mathcal{H}, D)$ and $a \in \mathcal{A}$,

$$\begin{aligned} [D^{-2z}, a] &= \left[\int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} d\lambda, a \right] \\ &= \int_{\gamma} \lambda^{-z} [(\lambda - D^2)^{-1}, a] d\lambda \\ &= \int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} [D^2, a] (\lambda - D^2)^{-1} d\lambda, \end{aligned}$$

then $[D^2, a] \notin B(\mathcal{H})$, so this is not a standard MOI.

Part 2: Abstract pseudodifferential calculus in the style of Connes–Moscovici, Higson, Guillemin

Pseudodifferential operators

On $\mathbb{R}^d$, a differential operator $L = \sum_{|\alpha| \leq k} a_\alpha(x) \partial^\alpha$ can be written as

$$L = \mathcal{F}^{-1} \circ M_{p_L} \circ \mathcal{F},$$

where M_{p_L} indicates multiplying with the polynomial $p_L(x, \xi) := \sum_{|\alpha| \leq k} a_\alpha(x) \xi^\alpha$.

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Generally speaking, a pseudodifferential operator of order k on $\mathbb{R}^d$ is an operator of the form $L = \mathcal{F}^{-1} \circ M_{p_L} \circ \mathcal{F}$ where the function p_L is more general, such that

$$L : \mathcal{H}^{s+k,2} \rightarrow \mathcal{H}^{s,2},$$

where

$$\mathcal{H}^{s,2}(\mathbb{R}^n) := \{f \in \mathcal{S}'(\mathbb{R}^n) : \mathcal{F}^{-1}[(1 + |\xi|^2)^{s/2} \mathcal{F}f] \in L_2(\mathbb{R}^n)\},$$

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On a smooth manifold M , we can define classes of pseudodifferential operators on $L_2(M)$ by patching together operators like this.

Sobolev spaces

Given an invertible, positive self-adjoint operator Θ on a separable Hilbert space $\mathcal{H}$, we can define the ‘Sobolev’ spaces $\mathcal{H}^s$, $s \in \mathbb{R}$, as the completion of $\text{dom } \Theta^s$ under the norm

$$\|\xi\|_s^2 = \langle \xi, \xi \rangle_s := \langle \Theta^s \xi, \Theta^s \xi \rangle_{\mathcal{H}} = \|\Theta^s \xi\|^2, \quad \xi \in \text{dom } \Theta^s.$$

This forms a Hilbert space. We have continuous embeddings

$$\mathcal{H}^t \subseteq \mathcal{H}^s, \quad s \leq t,$$

because

$$\|\Theta^s \xi\| \leq \|\Theta^{s-t}\|_{\infty} \|\Theta^t \xi\|.$$

We put

$$\mathcal{H}^{\infty} := \bigcap_{s \in \mathbb{R}} \mathcal{H}^s, \quad \mathcal{H}^{-\infty} := \bigcup_{s \in \mathbb{R}} \mathcal{H}^s,$$

and we get for free that $\mathcal{H}^{\infty}$ is dense in $\mathcal{H}$.

Analytic order

Even though Θ itself is an unbounded operator on $\mathcal{H}$, if we regard it as an operator

$$\Theta : \mathcal{H}^1 \rightarrow \mathcal{H}^0 = \mathcal{H},$$

it is a perfectly good bounded operator:

$$\|\Theta\|_{\mathcal{H}^1 \rightarrow \mathcal{H}^0} = \sup_{\xi: \|\Theta\xi\| \leq 1} \|\Theta\xi\| = 1.$$

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We define $\text{op}^r(\Theta)$ for $r \in \mathbb{R}$ as those $T : \mathcal{H}^\infty \rightarrow \mathcal{H}^\infty$ that extend to a bounded operator

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We define $\text{OP}^r(\Theta)$ as those $T \in \text{op}^r(\Theta)$ for which $\delta_\Theta^n(T) \in \text{op}^r(\Theta)$ for all $n \in \mathbb{N}$, where $\delta_\Theta(T) = [\Theta, T]$.

Examples

- If Δ is the Laplace operator on $\mathbb{R}^n$, setting $\Theta = (1 - \Delta)^{1/2}$ gives the standard (Bessel potential) Sobolev spaces. The k -th order (pseudo)differential operators are contained in $OP^k(\Theta)$.

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- Taking $\Theta = (1 - \Delta)^{1/2}$ where Δ is the sub-Laplacian on a stratified Lie group gives the Sobolev spaces defined by Folland and Stein.
- For a spectral triple $(\mathcal{A}, \mathcal{H}, D)$ it makes sense to put $\Theta = (1 + D^2)^{1/2}$. Then for example $D \in \text{OP}^1(\Theta)$, and for a *regular* spectral triple $a, [D, a] \in \text{OP}^0(\Theta)$ for all $a \in \mathcal{A}$.

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- If Θ is bounded, $\mathcal{H}^s \simeq \mathcal{H}$ and $\text{op}^r(\Theta) = B(\mathcal{H})$ for all $s, r \in \mathbb{R}$.

Elliptic operators

Our goal is to construct MOIs where all operators are in $\text{op}(\Theta)$. First, we need a functional calculus for such operators. Analogously to usual notions of pseudodifferential operators, a functional calculus can be constructed for *elliptic* operators.

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We define $T \in \text{op}^r(\Theta)$ to be Θ -*elliptic*, if there is a parametrix $P \in \text{op}^{-r}(\Theta)$ such that

$$TP = 1_{\mathcal{H}^\infty} + \text{op}^{-\infty}(\Theta);$$

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By a Borel Lemma argument, it suffices if

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Elliptic operators 2

For any spectral triple $(\mathcal{A}, \mathcal{H}, D)$ and $\Theta = (1 + D^2)^{1/2}$, we have that $D \in \text{op}^1(\Theta)$ is Θ -elliptic. Furthermore, $D + V$ is Θ -elliptic if $V \in \text{op}^r(\Theta)$ with $r < 1$.

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If $T \in \text{op}^r(\Theta)$ is Θ -elliptic,

- If $x \in \mathcal{H}^{-\infty}$, then $Tx \in \mathcal{H}^s$ implies that $x \in \mathcal{H}^{s+r}$ (*elliptic regularity*).
- If $T : \mathcal{H}^r \subseteq \mathcal{H}^0 \rightarrow \mathcal{H}^0$ (i.e. $r \geq 0$) is a symmetric operator, then it is self-adjoint given domain $\mathcal{H}^r$. This situation will be referred to as ' T is Θ -elliptic and symmetric'.

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$$T : \mathcal{H}^{r+s} \subseteq \mathcal{H}^s \rightarrow \mathcal{H}^s$$

is self-adjoint for any other $s \in \mathbb{R}$. In fact, these operators need not even be symmetric or normal.

Functional calculus

We write $f \in L^\beta_\infty(\mathbb{R})$ for some $\beta \in \mathbb{R}$ if $f(x)(1+x^2)^{-\beta/2} \in L_\infty(\mathbb{R})$.

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H.–McDonald–van Nuland (2024)

Let $T \in \text{op}^r(\Theta)$, $r > 0$, be Θ -elliptic and symmetric. If $f \in L_\infty^\beta(\mathbb{R})$, $\beta \in \mathbb{R}$, then

$$f(T) \in \text{op}^{r\beta}(\Theta).$$

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Furthermore, if A is self-adjoint on $\mathcal{H}$, $A \in \text{op}^t(\Theta)$, $t \in \mathbb{R}$, and A commutes strongly with T , then for $f \in L_\infty^\beta(\mathbb{R})$, $\beta \geq 0$, we have

$$f(A) \in \text{op}^{t\beta}(\Theta).$$

This second part applies for example to $i \frac{d}{dx}$ in $\text{op}(1 - \Delta)^{1/2}$ on $\mathbb{R}^d$.

Part 3: MOIs as pseudodifferential operators

Unbounded MOIs

H.–McDonald–van Nuland (2024)

Let $H_i \in \text{op}^{h_i}(\Theta)$, $h_i > 0$ Θ -elliptic and symmetric for $i = 0, \dots, n$, and $X_i \in \text{op}^{r_i}(\Theta)$ for $i = 1, \dots, n$. Let $\phi : \mathbb{R}^{n+1} \rightarrow \mathbb{C}$ such that

$$\phi(\lambda_0, \dots, \lambda_n) = \int_{\Omega} a_0(\lambda_0, \omega) \cdots a_n(\lambda_n, \omega) d\nu(\omega),$$

with finite measure space (Ω, ν) with $a_j(x, \omega)(1 + x^2)^{-\beta_j/2} : \mathbb{R} \times \Omega \rightarrow \mathbb{C}$ measurable and bounded. Then for $\psi \in \mathcal{H}^\infty$,

$$T_{\phi}^{H_0, \dots, H_n}(X_1, \dots, X_n)\psi := \int_{\Omega} a_0(H_0, \omega) X_1 a_1(H_1, \omega) \cdots X_n a_n(H_n, \omega) \psi d\nu(\omega)$$

is a well-defined vector in $\mathcal{H}^\infty$ independent of the representation of ϕ , and

$$T_{\phi}^{H_0, \dots, H_n} : \text{op}^{r_1}(\Theta) \times \cdots \times \text{op}^{r_n}(\Theta) \rightarrow \text{op}^{\sum_j r_j + \sum_j \beta_j h_j}(\Theta).$$

Unbounded MOIs: the useful bit

If $f \in C^{n+2}(\mathbb{R})$, and $f^{(k)} \in L_{\infty}^{\beta-k}(\mathbb{R})$ for $k = 0, \dots, n+2$, then for $H \in \text{op}^h(\Theta)$, $h > 0$ Θ -elliptic and symmetric, and $X_i \in \text{op}^{r_i}(\Theta)$,

$$T_{f^{[n]}}^{H, \dots, H}(X_1, \dots, X_n) \in \bigcap_{\varepsilon > 0} \text{op}^{(\beta-n)h + \sum_j r_j + \varepsilon}(\Theta).$$

If $f \in C^{\infty}(\mathbb{R})$ and $f^{(k)} \in L_{\infty}^{\beta-k}(\mathbb{R})$ for all $k \in \mathbb{N}$, we write $f \in S^{\beta}(\mathbb{R})$.

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For integral we saw earlier, with the appropriate conditions on D and a , we have

$$\int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} [D^2, a] (D^2 - \lambda)^{-1} d\lambda = T_{f^{[1]}}^{D^2, D^2}([D^2, a]) \in \bigcap_{\varepsilon > 0} \text{op}^{-2\Re(z)-1+\varepsilon}(\Theta),$$

with $f(x) = x^{-z}$ and $\Theta = (1 + D^2)^{\frac{1}{2}}$.

Two rules

MOIs as we defined them come with two identities:

- ① $f(A) - f(B) = T_{f^{[1]}}^{A,B}(A - B);$
- ② $[f(H), a] = T_{f^{[1]}}^{H,H}([H, a]),$

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$$\textcircled{2} \quad [f(H), a] = T_{f^{[1]}}^{H,H}([H, a]),$$

and the higher order analogues (since $T_{f^{[0]}}^H() = f(H)$)

$$\begin{aligned} \textcircled{1} \quad & T_{f^{[n]}}^{H_0, \dots, A, \dots, H_n}(V_1, \dots, V_n) - T_{f^{[n]}}^{H_0, \dots, B, \dots, H_n}(V_1, \dots, V_n) \\ &= T_{f^{[n+1]}}^{H_0, \dots, A, B, \dots, H_n}(V_1, \dots, A - B, \dots, V_n); \end{aligned}$$

$$\begin{aligned} \textcircled{2} \quad & T_{f^{[n]}}^{H_0, \dots, H_n}(V_1, \dots, V_{j-1}, aV_j, \dots, V_n) - T_{f^{[n]}}^{H_0, \dots, H_n}(V_1, \dots, V_{j-1}a, V_j, \dots, V_n) \\ &= T_{f^{[n+1]}}^{H_0, \dots, H_j, H_j, \dots, H_n}(V_1, \dots, V_{j-1}, [H_j, a], V_{j+1}, \dots, V_n). \end{aligned}$$

Taylor expansion

The first rule on its own gives a Taylor expansion:

$$\begin{aligned} f(H + V) &\stackrel{(1)}{=} f(H) + T_{f^{[1]}}^{H+V, H}(V) \\ &\stackrel{(1)}{=} f(H) + T_{f^{[1]}}^{H, H}(V) + T_{f^{[2]}}^{H+V, H, H}(V, V), \end{aligned}$$

and repeat. We get for all $N \in \mathbb{N}$

$$f(H + V) = \sum_{n=0}^N T_{f^{[n]}}^{H, \dots, H}(V, \dots, V) + T_{f^{[N+1]}}^{H+V, H, \dots, H}(V, \dots, V).$$

Note: if H and V commute,

$$T_{f^{[n]}}^{H, \dots, H}(V, \dots, V) = \frac{1}{n!} f^{(n)}(H) V^n.$$

Commutator expansion

In similar manner, by rule (2) we get

$$T_{f^{[n]}}^{H, \dots, H}(V_1, \dots, V_n) = V_1 T_{f^{[n]}}^{H, \dots, H}(1, V_2, \dots, V_n) + T_{f^{[n+1]}}^{H, \dots, H}([H, V], 1, V_2, \dots, V_n),$$

repeating and remembering that $T_{f^{[n]}}^{H, \dots, H}(1, \dots, 1) = \frac{1}{n!} f^{(n)}(H)$, we get

$$\begin{aligned} T_{f^{[n]}}^{H, \dots, H}(V_1, \dots, V_n) &= \sum_{m=0}^N \sum_{m_1 + \dots + m_n = m} \frac{C_{m_1, \dots, m_n}}{(n+m)!} \delta_H^{m_1}(X_1) \cdots \delta_H^{m_n}(X_n) f^{(n+m)}(H) \\ &\quad + S_{H, V}^N. \end{aligned}$$

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The combinatorics to get this expression is exactly the same as how one gets the cocycle of the local index formula, writing $A^{(k)} := \delta_{D^2}^n(A)$,

$$\begin{aligned} &\phi_n(a_0, \dots, a_n) \\ &= \sum_{|k|, q \geq 0} c_{n, k, q} \operatorname{Res}_{z=0} z^q \operatorname{Tr} \left(a_0 [D, a_1]^{(k_1)} \cdots [D, a_n]^{(k_n)} |D|^{-2|k| - 2z - n} \right), \end{aligned}$$

Asymptotic expansions

We say that $T \sim \sum_{k=0}^{\infty} T_k$ for $T, T_k \in \text{op}(\Theta)$ if

$$T - \sum_{k=1}^N T_k \in \text{op}^{m_N}(\Theta), \quad m_N \downarrow -\infty.$$

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If $f \in S^{\beta}(\mathbb{R})$, if $H \in \text{op}^h(\Theta)$, $h > 0$ is Θ -elliptic and symmetric, and if $V \in \text{op}^r(\Theta)$ with $r < h$, then

$$f(H + V) \sim \sum_{n=0}^{\infty} T_{f^{[n]}}^{H, \dots, H}(V, \dots, V).$$

If furthermore $\delta_H^n(V) \in \text{op}^{r+n(h-\varepsilon)}(\Theta)$ for some $\varepsilon > 0$ (for example $H = \Theta$, $V \in \text{OP}^r$)

$$T_{f^{[n]}}^{H, \dots, H}(V, \dots, V) \sim \sum_{m=0}^{\infty} \sum_{m_1 + \dots + m_n = m} \frac{C_{m_1, \dots, m_n}}{(n+m)!} \delta_H^{m_1}(V) \dots \delta_H^{m_n}(V) f^{(n+m)}(H).$$

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Combined,

$$f(H + V) \sim \sum_{n, m=0}^{\infty} \sum_{m_1 + \dots + m_n = m} \frac{C_{m_1, \dots, m_n}}{(n+m)!} \delta_H^{m_1}(V) \dots \delta_H^{m_n}(V) f^{(n+m)}(H).$$

A familiar expansion

Recall that

$$[f(\Theta), X] = T_{f^{[1]}}^{\Theta, \Theta}([\Theta, X]).$$

Therefore, for $X \in \text{OP}^r(\Theta)$, the expansions on the last slide give

$$[f(\Theta), X] \sim \sum_{k=1}^{\infty} \frac{1}{k!} \delta_{\Theta}^k(X) f^{(k)}(\Theta).$$

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In particular,

$$[\Theta^{\alpha}, X] \sim \sum_{k=1}^{\infty} \binom{\alpha}{k} \delta_{\Theta}^k(X) \Theta^{\alpha-k}, \quad \alpha \in \mathbb{C},$$

and

$$[\log(\Theta), X] \sim \sum_{k=1}^{\infty} \frac{(-1)^{k-1}}{k} \delta_{\Theta}^k(X) \Theta^{-k},$$

and we have that $[\Theta^{\alpha}, X] \in \text{OP}^{r+\Re(\alpha)-1}(\Theta)$ and $[\log(\Theta), X] \in \text{OP}^{r-1}(\Theta)$.

Asymptotic trace expansions

H.–McDonald–van Nuland (2024)

Let $(\mathcal{A}, \mathcal{H}, D)$ be a regular s -summable spectral triple $((1 + D^2)^{-\frac{1}{2}} \in \mathcal{L}_s)$. Let V self-adjoint and bounded, generated by $\mathcal{A}$ and D . Then if f has sufficient regularity, as $t \rightarrow 0$,

$$\begin{aligned} & \mathrm{Tr}(f(tD + tV)) \\ &= \sum_{n=0}^N \sum_{m=0}^N \sum_{m_1 + \dots + m_n = m} t^{n+m} \frac{C_{m_1, \dots, m_n}}{(n+m)!} \mathrm{Tr}(\delta_D^{m_1}(V) \cdots \delta_D^{m_n}(V) f^{(n+m)}(tD)) \\ &+ O(t^{N+1-s}). \end{aligned}$$

Bonus: Functional calculus for OP

If $A \in \text{op}^r(\Theta)$ is Θ -elliptic and symmetric, then by rule (2) we know that for $f \in L^\beta_\infty(\mathbb{R})$

$$f(A) \in \text{op}^{\beta r}(\Theta),$$

$$[\Theta, f(A)] = T_{f[1]}^{A,A}([\Theta, A]),$$

and similar expressions hold for $\delta_\Theta^n(f(A))$. If $A \in \text{OP}^r(\Theta)$, then we can deduce what the order is of these expressions if $f \in S^\beta(\mathbb{R})$, so that

$$\delta_\Theta^n(f(A)) \in \bigcap_{\varepsilon > 0} \text{op}^{r\beta + \varepsilon}(\Theta).$$

We therefore conclude that $f(A) \in \bigcap_{\varepsilon > 0} \text{OP}^{r\beta + \varepsilon}(\Theta)$.

Thanks

Thank you for your attention!