The Joy of MOIs Cortona 2024

Eva-Maria Hekkelman

UNSW

June 26 2024

EI	М.	Hek	kelman	- (UN	JSW)

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Summary of this talk

- Motivation for studying MOIs
- Pseudodifferential calculus
- MOIs of pseudodifferential operators

This talk is based on joint work with Ed McDonald and Teun van Nuland.

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Part 1: Motivation

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Exhibit A

We use the Chern character of $(\mathcal{A}, \mathcal{H}, D)$ in entire cyclic cohomology (cf. [2]) given in the most efficient manner by the JLO formula, which defines the components of an entire cocycle in the (b, B) bicomplex:

(90)
$$\psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0, \dots, a^n) = \sqrt{2i} \int_{\sum_{i=0}^n v_i = 1, v_i \ge 0}^n \psi_n(a^0,$$

Trace
$$\left(a^{0} e^{-v_{0}D^{2}} [D, a^{1}] e^{-v_{1}D^{2}} \dots e^{-v_{n-1}D^{2}} [D, a^{n}] e^{-v_{n}D^{2}}\right)$$
, $\forall a^{j} \in \mathcal{A}$

where n is odd.

We introduce a parameter ϵ by replacing D^2 by ϵD^2 , which yields a cocycle ψ_n^{ϵ} which is cohomologous to ψ_n . One has moreover

(91)
$$\psi_n^{\epsilon}(a^0,\ldots,a^n) = \sqrt{2i} \left(\int_{\sum_{0}^n v_t=1}^n \theta(\epsilon \ v_0,\ldots,\epsilon \ v_n) \ \pi \ dv_i \right) \ \epsilon^{n/2} ,$$

From [ConnesMoscovici1995]

Exhibit B

Let us now show that if $b \in \cap$ Dom $L^k \mathbb{R}^q$ then $b \in$ Dom δ . The proof is more subtle than one would expect, because the obvious argument, using

$$|D| = \pi^{-1} \int_0^\infty \frac{D^2}{D^2 + \mu} \ \mu^{-1/2} \ d\mu \ ,$$

requires some care. Indeed, one gets from the above

$$[|D|,b] = \pi^{-1} \int_0^\infty (D^2 + \mu)^{-1} \ [D^2,b] \ (D^2 + \mu)^{-1} \ \mu^{1/2} \ d\mu \ .$$

We can replace $[D^2, b]$ by |D|, which has the same size, and get

$$\int_0^\infty (D^2 + \mu)^{-2} |D| \ \mu^{1/2} \ d\mu = \int_0^\infty (1 + t)^{-2} \ t^{1/2} \ dt \ .$$

For this to work, we need to move $[D^2,b]$ in front of the above integral, i.e. use the finiteness of the norm of

$$\int_0^\infty \underbrace{[(D^2+\mu)^{-1}, [D^2, b]]}_{-(D^2+\mu)^{-1}[D^2, [D^2, b]](D^2+\mu)^{-1}} (D^2+\mu)^{-1} \ \mu^{1/2} \ d\mu \ .$$

This finiteness follows from:

1)
$$(D^2 + \mu)^{-1} [D^2, [D^2, b]]$$
 bounded since $b \in \text{Dom} L^2$

2)
$$\int_0^\infty \|(D^2 + \mu)^{-2}\| \mu^{1/2} d\mu \le C \int_0^1 \mu^{1/2} d\mu + \int_1^\infty \mu^{-3/2} d\mu < \infty.$$

Once $[D^2, b]$ is moved in front the above calculation applies.

From [ConnesMoscovici1995]

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Exhibit C

Now, onwards with the computation, the first part of which is straightforward:

$$\begin{split} [\Delta^{-z}, A] B &= \frac{1}{2\pi i} \int \lambda^{-z} [(\lambda - \Delta)^{-1}, A] B \, d\lambda \\ &= \frac{1}{2\pi i} \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] (\lambda - \Delta)^{-1} B \, d\lambda \\ &= \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] B (\lambda - \Delta)^{-1} \, d\lambda \\ &+ \int \lambda^{-z} (\lambda - \Delta)^{-1} [\Delta, A] (\lambda - \Delta)^{-1} [\Delta, B] (\lambda - \Delta)^{-1} \, d\lambda \end{split}$$

(In the last step we did two things at once: we commuted B past $(\lambda - \Delta)^{-1}$ and we then used the formula $[S^{-1}, T] = S^{-1}[T, S]S^{-1}$.) The operators $[\Delta, A]$ and $[\Delta, B]$ have orders 1 and 2, respectively.

Before going on, we shall introduce some better notation for our contour integrals.

2.5 Definition. If $D_0, ..., D_p$ are differential operators on the closed manifold M, then denote by $I_z(D_0, ..., D_p)$ the integral

$$\frac{1}{2\pi i}\int \lambda^{-z} D_0(\lambda-\Delta)^{-1}\cdots D_p(\lambda-\Delta)^{-1} d\lambda$$

(in the integral, copies of $(\lambda - \Delta)^{-1}$ alternate with the operators D_j). The integral converges if Re(z) < n, in the sense we discussed above, and defines an operator on $C^{\infty}(M)$.

From [Higson2003]

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Exhibit D

Theorem 4.2 (Semifinite Odd Local Index Theorem). Let $(\mathcal{A}, \mathcal{H}, \mathcal{D})$ be an odd finitely summable QC^{∞} spectral triple with spectral dimension $p \geq 1$. Let $N = \lfloor p/2 \rfloor + 1$ where $\lfloor \cdot \rfloor$ denotes the integer part, and let $u \in \mathcal{A}$ be unitary. Then

1)
$$sf(\mathcal{D}, u^*\mathcal{D}u) = \frac{1}{\sqrt{2\pi i}} res_{r=(1-p)/2} \left(\sum_{m=1,odd}^{2N-1} \phi_m^r(Ch_m(u)) \right)$$

where for $a_0, ..., a_m \in \mathcal{A}$, $l = \{a + iv : v \in \mathbf{R}\}$, 0 < a < 1/2, $R_s(\lambda) = (\lambda - (1 + s^2 + \mathcal{D}^2))^{-1}$ and r > 0 we define $\phi_m^r(a_0, a_1, ..., a_m)$ to be

$$\frac{-2\sqrt{2\pi i}}{\Gamma((m+1)/2)}\int_0^\infty s^m \tau\left(\frac{1}{2\pi i}\int_l^{\lambda^{-p/2-r}}a_0R_s(\lambda)[\mathcal{D},a_1]R_s(\lambda)\cdots[\mathcal{D},a_m]R_s(\lambda)d\lambda\right)ds.$$

In particular the sum on the right hand side of 1) analytically continues to a deleted neighbourhood of r = (1 - p)/2 with at worst a simple pole at r = (1 - p)/2. Moreover, the complex function-valued cochain $(\phi_m^r)_{m=1,odd}^{2N-1}$ is a (b, B) cocycle for \mathcal{A} modulo functions holomorphic in a half-plane containing r = (1 - p)/2.

From [CareyPhillipsRennieSukochev2006]

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Exhibit E

Let us introduce the following convenient notation (cf. [10]). If A_0, \ldots, A_n are operators, we define a *t*-dependent quantity by

$$\langle A_0, \dots, A_n \rangle_n := t^n \operatorname{Tr} \int_{\Delta_n} A_0 e^{-s_0 t D^2} A_1 e^{-s_1 t D^2} \cdots A_n e^{-s_n t D^2} d^n s.$$
(3)

Note the difference in notation with [10], for which the same symbol is used for the supertrace of the same expression, rather than the trace. Also, we are integrating over the 'inflated' *n*-simplex $t\Delta^n$, yielding the factor t^n . The forms $\langle A_0, \ldots, A_n \rangle$ satisfy, *mutatis mutandis*, the following properties.

Lemma 7. (See [10].) In each of the following cases, we assume that the operators A_i are such that each term is well defined:

1.
$$\langle A_0, ..., A_n \rangle_n = \langle A_i, ..., A_n, ..., A_{i-1} \rangle_n;$$

2. $\langle A_0, ..., A_n \rangle_n = \sum_{i=0}^n \langle 1, ..., A_i, ..., A_n, A_0, ..., A_{i-1} \rangle_n;$
3. $\sum_{i=0}^n \langle A_0, ..., [D, A_i], ..., A_n \rangle_n = 0;$
4. $\langle A_0, ..., [D^2, A_i], ..., A_n \rangle_n = \langle A_0, ..., A_{i-1}A_i, ..., A_n \rangle_{n-1} - \langle A_0, ..., A_iA_{i+1}, ..., A_n \rangle_{n-1}.$

From [vanSuijlekom2011]

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Multiple operator integrals

Let $\phi : \mathbb{R}^{n+1} \to \mathbb{C}$ be such that

$$\phi(\lambda_0,\ldots,\lambda_n)=\int_\Omega a_0(\lambda_0,\omega)\cdots a_n(\lambda_n,\omega)d
u(\omega),$$

with finite measure space (Ω, ν) and measurable and bounded $a_j : \mathbb{R} \times \Omega \to \mathbb{C}$.

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with finite measure space (Ω, ν) and measurable and bounded $a_j : \mathbb{R} \times \Omega \to \mathbb{C}$.

Let H_0, \ldots, H_n be self-adjoint, for $V_1, \ldots, V_n \in B(\mathcal{H})$ define the MOI

$$T_{\phi}^{H_0,\ldots,H_n}(V_1,\ldots,V_n)\psi$$

:= $\int_{\Omega} a_0(H_0,\omega)V_1a_1(H_1,\omega)\cdots V_na_n(H_n,\omega)\psi d\nu(\omega), \quad \psi \in \mathcal{H}.$

Then,

$$T_{\phi}^{H_0,...,H_n}:B(\mathcal{H}) imes\cdots imes B(\mathcal{H}) o B(\mathcal{H})$$

and this does not depend on how we represent ϕ (its *symbol*).

E.-M. Hekkelman (UNSW)

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Divided differences

Symbols of MOIs encountered in the wild are almost always divided differences, which are defined recursively for $f \in C^n(\mathbb{R})$ as

$$f^{[0]}(\lambda) := f(\lambda);$$

$$f^{[n]}(\lambda_0, \dots, \lambda_n) := \frac{f^{[n-1]}(\lambda_0, \dots, \lambda_{n-1}) - f^{[n-1]}(\lambda_1, \dots, \lambda_n)}{\lambda_0 - \lambda_n},$$

with an appropriate limit if $\lambda_0 = \lambda_n$. In particular,

$$\frac{1}{n!}f^{(n)}(\lambda)=f^{[n]}(\lambda,\ldots,\lambda).$$

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Example MOIs

For example the JLO cocycle is

$$\int_{\Delta_n} \operatorname{Tr}(\eta a_0 e^{-t_0 D^2}[D, a_1] e^{-t_1 D^2} \cdots [D, a_n] e^{-t_n D^2}) dt$$

= $\operatorname{Tr}(\eta a_0 T_{f^{[n]}}^{D^2}([D, a_1], \dots, [D, a_n])),$

with $f(x) = \exp(-x)$.

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$$\frac{d^n}{dt^n}f(H+tV)|_{t=0} = T^{H,...,H}_{f^{[n]}}(V,...,V),$$

each of which has been used to obtain sharp estimates. (Potapov, Sukochev, Skripka, Caspers, Montgomery-Smith, McDonald, Peller, ...)

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Furthermore, MOIs can systematise operator integral techniques in NCG.

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A problem

If you write, like in *The Local Index Formula in Noncommutative Geometry* by Nigel Higson, for a spectral triple $(\mathcal{A}, \mathcal{H}, D)$ and $a \in \mathcal{A}$,

$$\begin{split} [D^{-2z},a] &= \left[\int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} d\lambda, a\right] \\ &= \int_{\gamma} \lambda^{-z} [(\lambda - D^2)^{-1}, a] d\lambda \\ &= \int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} [D^2, a] (\lambda - D^2)^{-1} d\lambda, \end{split}$$

then $[D^2, a] \notin B(\mathcal{H})$, so this is not a standard MOI.

Part 2: Abstract pseudodifferential calculus in the style of Connes–Moscovici, Higson, Guillemin

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Pseudodifferential operators

On \mathbb{R}^d , a differential operator $L = \sum_{|\alpha| \leq k} a_{\alpha}(x) \partial^{\alpha}$ can be written as

$$L = \mathcal{F}^{-1} \circ M_{p_L} \circ \mathcal{F},$$

where M_{p_L} indicates multiplying with the polynomial $p_L(x,\xi) := \sum_{|\alpha| \le k} a_{\alpha}(x)\xi^{\alpha}$.

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Generally speaking, a pseudodifferential operator of order k on \mathbb{R}^d is an operator of the form $L = \mathcal{F}^{-1} \circ M_{p_L} \circ \mathcal{F}$ where the function p_L is more general, such that

$$L: \mathcal{H}^{s+k,2} \to \mathcal{H}^{s,2},$$

where

$$\mathcal{H}^{s,2}(\mathbb{R}^n):=\{f\in\mathcal{S}'(\mathbb{R}^n):\mathcal{F}^{-1}ig[(1+|\xi|^2)^{s/2}\mathcal{F}fig]\in L_2(\mathbb{R}^n)\},$$

are Bessel potential Sobolev spaces.

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are Bessel potential Sobolev spaces.

On a Riemannian manifold M, we can define classes of pseudodifferential operators on $L_2(M)$ that locally look like above.

Sobolev spaces

Given an invertible, positive self-adjoint operator Θ on a separable Hilbert space \mathcal{H} , we can define the 'Sobolev' spaces \mathcal{H}^s , $s \in \mathbb{R}$, as the completion of dom Θ^s under the norm

$$\|\xi\|_s^2 = \langle \xi, \xi \rangle_s := \langle \Theta^s \xi, \Theta^s \xi \rangle_{\mathcal{H}} = \|\Theta^s \xi\|^2, \quad \xi \in \operatorname{dom} \Theta^s$$

This forms a Hilbert space. We have continuous embeddings

$$\mathcal{H}^t \subseteq \mathcal{H}^s, \quad s \leq t,$$

because

$$\|\Theta^{s}\xi\| \leq \|\Theta^{s-t}\|_{\infty} \|\Theta^{t}\xi\|.$$

We put

$$\mathcal{H}^{\infty} := \bigcap_{s \in \mathbb{R}} \mathcal{H}^{s}, \quad \mathcal{H}^{-\infty} := \bigcup_{s \in \mathbb{R}} \mathcal{H}^{s},$$

and we get for free that \mathcal{H}^∞ is dense in $\mathcal{H}.$

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Analytic order

Even though Θ itself is an unbounded operator on $\mathcal H,$ if we regard it as an operator

$$\Theta: \mathcal{H}^1 \to \mathcal{H}^0 = \mathcal{H},$$

it is a perfectly good bounded operator:

$$\|\Theta\|_{\mathcal{H}^1
ightarrow \mathcal{H}^0} = \sup_{\xi: \|\Theta\xi\| \leq 1} \|\Theta\xi\| = 1.$$

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We define op^r(Θ) for $r \in \mathbb{R}$ as those $T : \mathcal{H}^{\infty} \to \mathcal{H}^{\infty}$ that extend to a bounded operator

$$T: \mathcal{H}^{s+r} \to \mathcal{H}^s, \quad s \in \mathbb{R}.$$

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$$\mathcal{T}:\mathcal{H}^{s+r}
ightarrow\mathcal{H}^{s},\quad s\in\mathbb{R}.$$

We define $OP'(\Theta)$ as those $T \in op^r(\Theta)$ for which $[\Theta, T] \in op^r(\Theta)$, $[\Theta, [\Theta, T]] \in op^r(\Theta), \ \delta^n_{\Theta}(T) \in op^r(\Theta)$.

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If Δ is the Laplace operator on ℝⁿ, setting Θ = (1 − Δ)^{1/2} gives the standard (Bessel potential) Sobolev spaces. The k-th order (pseudo)differential operators are contained in OP^k(Θ).

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- Taking $\Theta = (1 \Delta)^{1/2}$ where Δ is the sub-Laplacian on a stratified Lie group gives the Sobolev spaces defined by Folland and Stein.

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- Taking $\Theta = (1 \Delta)^{1/2}$ where Δ is the sub-Laplacian on a stratified Lie group gives the Sobolev spaces defined by Folland and Stein.
- For a spectral triple $(\mathcal{A}, \mathcal{H}, D)$ it makes sense to put $\Theta = (1 + D^2)^{1/2}$. Then for example $D \in OP^1(\Theta)$, and for a *regular* spectral triple $a, [D, a] \in OP^0(\Theta)$ for all $a \in \mathcal{A}$.

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- Taking $\Theta = (1 \Delta)^{1/2}$ where Δ is the sub-Laplacian on a stratified Lie group gives the Sobolev spaces defined by Folland and Stein.
- For a spectral triple (A, H, D) it makes sense to put Θ = (1 + D²)^{1/2}. Then for example D ∈ OP¹(Θ), and for a *regular* spectral triple a, [D, a] ∈ OP⁰(Θ) for all a ∈ A.
- If Θ is bounded, $\mathcal{H}^{s} \simeq \mathcal{H}$ and $\operatorname{op}^{r}(\Theta) = B(\mathcal{H})$ for all $s, r \in \mathbb{R}$.

Our goal is to construct MOIs where all operators are in $op(\Theta)$. First, we need a functional calculus for such operators. Analogously to usual notions of pseudodifferential operators, a functional calculus can be constructed for *elliptic* operators.

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We define $T \in op^{r}(\Theta)$ to be *elliptic*, if there is a parametrix $P \in op^{-r}(\Theta)$ such that

$$TP = 1_{\mathcal{H}^{\infty}} + op^{-\infty}(\Theta);$$

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By a Borel Lemma argument, it suffices if

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For any spectral triple $(\mathcal{A}, \mathcal{H}, D)$ and $\Theta = (1 + D^2)^{1/2}$, we have that $D \in op^1(\Theta)$ is elliptic. Furthermore, D + V is elliptic if $V \in op^r$ with r < 1.

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If $T \in \mathsf{op}^r(\Theta)$ is elliptic,

- If $x \in \mathcal{H}^{-\infty}$, then $Tx \in \mathcal{H}^s$ implies that $x \in \mathcal{H}^{s+r}$ (elliptic regularity).
- If T : H^r ⊆ H⁰ → H⁰ (i.e. r ≥ 0) is a symmetric operator, then it is self-adjoint. This situation will be referred to as 'T is elliptic and symmetric'.

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$$T:\mathcal{H}^{r+s}\subseteq\mathcal{H}^s
ightarrow\mathcal{H}^s$$

is self-adjoint for any other $s \in \mathbb{R}$. In fact, these operators need not even be symmetric or normal.

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Functional calculus

We write $f \in L^{\beta}_{\infty}(\mathbb{R})$ for some $\beta \in \mathbb{R}$ if $f(x)(1+x^2)^{-\beta/2} \in L_{\infty}(\mathbb{R})$.

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H.-McDonald-van Nuland (2024)

Let $T \in op^{r}(\Theta)$, r > 0, be elliptic and symmetric. If $f \in L^{\beta}_{\infty}(\mathbb{R})$, then

 $f(T) \in op^{r\beta}(\Theta).$

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Furthermore, if A is self-adjoint on \mathcal{H} , $A \in op^t(\Theta)$, $t \in \mathbb{R}$, and A commutes strongly with T, then for $f \in L^{\beta}_{\infty}(\mathbb{R})$, $\beta \geq 0$, we have

 $f(A) \in op^{t\beta}(\Theta).$

This second part applies for example to $i\frac{d}{dx}$ in $op(1-\Delta)^{1/2}$ on \mathbb{R}^d .

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Part 3: MOIs as pseudodifferential operators

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Unbounded MOIs

H.-McDonald-van Nuland (2024)

Let $H_i \in op^{h_i}(\Theta)$, $h_i > 0$ elliptic and symmetric for i = 0, ..., n, and $X_i \in op^{r_i}(\Theta)$ for i = 1, ..., n. Let $\phi : \mathbb{R}^{n+1} \to \mathbb{C}$ such that

$$\phi(\lambda_0,\ldots,\lambda_n)=\int_\Omega a_0(\lambda_0,\omega)\cdots a_n(\lambda_n,\omega)d
u(\omega),$$

with finite measure space (Ω, ν) with $a_j(x, \omega)(1 + x^2)^{-\beta_j/2} : \mathbb{R} \times \Omega \to \mathbb{C}$ measurable and bounded. Then for $\psi \in \mathcal{H}^{\infty}$,

$$T_{\phi}^{H_0,\ldots,H_n}(X_1,\ldots,X_n)\psi := \int_{\Omega} a_0(H_0,\omega)X_1a_1(H_1,\omega)\cdots X_na_n(H_n,\omega)\psi d
u(\omega)$$

is a well-defined vector in \mathcal{H}^∞ independent of the representation of $\phi,$ and

$$T^{H_0,\ldots,H_n}_{\phi}: \operatorname{op}^{r_1}(\Theta) \times \cdots \times \operatorname{op}^{r_n}(\Theta) \to \operatorname{op}^{\sum_j r_j + \sum_j \beta_j h_j}(\Theta).$$

Unbounded MOIs: the useful bit

If $f \in C^{n+2}(\mathbb{R})$, and $f^{(k)} \in L^{\beta-k}_{\infty}(\mathbb{R})$ for k = 0, ..., n+2, then for $H \in \operatorname{op}^{h}(\Theta)$, h > 0 elliptic and symmetric, and $X_{i} \in \operatorname{op}^{r_{i}}(\Theta)$,

$$T_{f^{[n]}}^{H,\ldots,H}(X_1,\ldots,X_n)\in\bigcap_{\varepsilon>0}\operatorname{op}^{(\beta-n)h+\sum_j r_j+\varepsilon}(\Theta).$$

If $f \in C^{\infty}(\mathbb{R})$ and $f^{(k)} \in L^{\beta-k}_{\infty}(\mathbb{R})$ for all $k \in \mathbb{N}$, we write $f \in S^{\beta}(\mathbb{R})$.

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If $f \in C^{n+2}(\mathbb{R})$, and $f^{(k)} \in L^{\beta-k}_{\infty}(\mathbb{R})$ for k = 0, ..., n+2, then for $H \in op^{h}(\Theta)$, h > 0 elliptic and symmetric, and $X_{i} \in op^{r_{i}}(\Theta)$,

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If $f \in C^{\infty}(\mathbb{R})$ and $f^{(k)} \in L^{\beta-k}_{\infty}(\mathbb{R})$ for all $k \in \mathbb{N}$, we write $f \in S^{\beta}(\mathbb{R})$. The integral we saw earlier,

$$\int_{\gamma} \lambda^{-z} (\lambda - D^2)^{-1} [D^2, a] (D^2 - \lambda)^{-1} d\lambda = T_{f^{[1]}}^{D^2, D^2} ([D^2, a]),$$

with $f(x) = x^{-z}$.

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Two rules

MOIs as we defined them come with two identities:

•
$$f(A) - f(B) = T_{f^{[1]}}^{A,B}(A - B);$$

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$$[f(H), a] = T_{f^{[1]}}^{H,H}([H, a]),$$

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$$[f(H), a] = T_{f^{[1]}}^{n, n}([H, a]),$$

and the higher order analogues (since $T_{f^{[0]}}^H() = f(H)$)

•
$$T_{f^{[n]}}^{H_0,...,A,...,H_n}(V_1,...,V_n) - T_{f^{[n]}}^{H_0,...,B,...,H_n}(V_1,...,V_n)$$

= $T_{f^{[n+1]}}^{H_0,...,A,B,...,H_n}(V_1,...,A-B,...,V_n);$

$$T_{f^{[n]}}^{H_0,\ldots,H_n}(V_1,\ldots,V_{j-1},aV_j,\ldots,V_n) - T_{f^{[n]}}^{H_0,\ldots,H_n}(V_1,\ldots,V_{j-1}a,V_j,\ldots,V_n) = T_{f^{[n+1]}}^{H_0,\ldots,H_j,H_j,\ldots,H_n}(V_1,\ldots,V_{j-1},[H_j,a],V_{j+1},\ldots,V_n).$$

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Taylor expansion

The first rule on its own gives a Taylor expansion:

$$f(H+V) \stackrel{(1)}{=} f(H) + T_{f^{[1]}}^{H+V,H}(V)$$
$$\stackrel{(1)}{=} f(H) + T_{f^{[1]}}^{H,H}(V) + T_{f^{[2]}}^{H+V,H,H}(V,V),$$

and repeat. We get for all $N \in \mathbb{N}$

$$f(H+V) = \sum_{n=0}^{N} T_{f^{[n]}}^{H,\ldots,H}(V,\ldots,V) + T_{f^{[N+1]}}^{H+V,H,\ldots,H}(V,\ldots,V).$$

Note: if H and V commute,

$$T_{f^{[n]}}^{H,...,H}(V,...,V) = \frac{1}{n!}f^{(n)}(H)V^{n}.$$

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Commutator expansion

In similar manner, by rule (2) we get

$$T_{f^{[n]}}^{H,\ldots,H}(V_1,\ldots,V_n) = V_1 T_{f^{[n]}}^{H,\ldots,H}(1,V_2,\ldots,V_n) + T_{f^{[n+1]}}^{H,\ldots,H}([H,V],1,V_2,\ldots,V_n),$$

repeating and remembering that $T^{H,\ldots,H}_{f^{[n]}}(1,\ldots,1)=rac{1}{n!}f^{(n)}(H)$, we get

$$T_{f^{[n]}}^{H,...,H}(V_1,...,V_n) = \sum_{m=0}^{N} \sum_{\substack{m_1+\cdots+m_n=m}} \frac{C_{m_1,...,m_n}}{(n+m)!} \delta_H^{m_1}(X_1) \cdots \delta_H^{m_n}(X_n) f^{(n+m)}(H) + S_{H,V}^N.$$

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The combinatorics to get this expression is exactly the same as how one gets the cocycle of the local index formula, writing $A^{(k)} := \delta_{D^2}^n(A)$,

$$\phi_n(a_0,\ldots,a_n) = \sum_{|k|,q\geq 0} c_{n,k,q} \operatorname{Res}_{z=0} z^q \operatorname{Tr}\left(a_0[D,a_1]^{(k_1)}\cdots[D,a_n]^{(k_n)}|D|^{-2|k|-2z-n}\right),$$

Asymptotic expansions

We say that $T \sim \sum_{k=0}^{\infty} T_k$ for $T, T_k \in op(\Theta)$ if $T - \sum_{k=1}^{N} T_k \in op^{m_N}(\Theta), \quad m_N \downarrow -\infty.$

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Asymptotic expansions

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m op}(\Theta)$ if

$$T-\sum_{k=1}^N T_k\in {
m op}^{m_N}(\Theta), \quad m_N\downarrow -\infty.$$

If $f \in S^{\beta}(\mathbb{R})$, if $H \in op^{h}(\Theta)$, h > 0 is elliptic and symmetric, and if $V \in op^{r}(\Theta)$ with r < h, then

$$f(H+V)\sim \sum_{n=0}^{\infty}T_{f^{[n]}}^{H,\ldots,H}(V,\ldots,V).$$

If furthermore $\delta^n_H(V) \in {
m op}^{r+n(h-arepsilon)}$ for some arepsilon>0 (for example $H=\Theta$, $V\in {
m OP}^r$)

$$T_{f^{[n]}}^{H,...,H}(V,...,V) \sim \sum_{m=0}^{\infty} \sum_{m_1+\dots+m_n=m} \frac{C_{m_1,...,m_n}}{(n+m)!} \delta_H^{m_1}(V) \cdots \delta_H^{m_n}(V) f^{(n+m)}(H).$$

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Combined,

$$f(H+V) \sim \sum_{n,m=0}^{\infty} \sum_{m_1+\dots+m_n=m} \frac{C_{m_1,\dots,m_n}}{(n+m)!} \delta_H^{m_1}(V) \cdots \delta_H^{m_n}(V) f^{(n+m)}(H).$$

E.-M. Hekkelman (UNSW)

The Joy of MOIs

A familiar expansion

Recall that

$$[f(\Theta), X] = T_{f^{[1]}}^{\Theta, \Theta}([\Theta, X]).$$

Therefore, for $X \in OP^{r}(\Theta)$, the expansions on the last slide give

$$[f(\Theta), X] \sim \sum_{k=1}^{\infty} \frac{1}{k!} \delta_{\Theta}^k(X) f^{(k)}(\Theta).$$

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$$[f(\Theta), X] \sim \sum_{k=1}^{\infty} \frac{1}{k!} \delta_{\Theta}^k(X) f^{(k)}(\Theta).$$

In particular,

$$[\Theta^{\alpha}, X] \sim \sum_{k=1}^{\infty} \binom{\alpha}{k} \delta^{k}_{\Theta}(X) \Theta^{\alpha-k}, \quad \alpha \in \mathbb{C},$$

and

$$[\log(\Theta), X] \sim \sum_{k=1}^{\infty} \frac{(-1)^{k-1}}{k} \delta_{\Theta}^k(X) \Theta^{-k},$$

and we have that $[\Theta^{\alpha}, X] \in OP^{r+\Re(\alpha)-1}(\Theta)$ and $[\log(\Theta), X] \in OP^{r-1}(\Theta)$.

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Asymptotic trace expansions

H.-McDonald-van Nuland (2024)

Let $(\mathcal{A}, \mathcal{H}, D)$ be a regular s-summable spectral triple $((1 + D^2)^{-1/2} \in \mathcal{L}_s)$. Let V self-adjoint and bounded, generated by \mathcal{A} and D. Then as $t \to 0$,

$$\operatorname{Tr}(f(tD + tV))$$

$$= \sum_{n=0}^{N} \sum_{m=0}^{N} \sum_{m_1 + \dots + m_n = m}^{N} t^{n+m} \frac{C_{m_1,\dots,m_n}}{(n+m)!} \operatorname{Tr}(\delta_D^{m_1}(V) \cdots \delta_D^{m_n}(V) f^{(n+m)}(tD))$$

$$+ O(t^{N+1-s}).$$

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Bonus: Functional calculus for OP

If $A \in \operatorname{op}^{r}(\Theta)$ is elliptic and symmetric, then by rule (2) we know that for $f \in L^{\beta}_{\infty}(\mathbb{R})$ $f(A) \in \operatorname{op}^{\beta r}(\Theta).$

$$[\Theta, f(A)] = T^{A,A}_{f^{[1]}}([\Theta, A]),$$

and similar expressions hold for $\delta^n_{\Theta}(f(A))$. If $A \in OP^r(\Theta)$, then we can deduce what the order is of these expressions if $f \in S^{\beta}(\mathbb{R})$, so that

$$\delta^n_\Theta(f(\mathcal{A}))\in igcap_{arepsilon>0} {
m op}^{reta+arepsilon}(\Theta).$$

We therefore conclude that $f(A) \in \bigcap_{\varepsilon > 0} \mathsf{OP}^{r\beta + \varepsilon}(\Theta)$.

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Thanks

Thank you for your attention!

EM.	Hekkelma	n (UNSW)
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